

Dynamic resistance measurement in a four-tape YBCO stack with various applied field orientation

Yanchao Liu, Zhenan Jiang, Quan Li, Chris W. Bumby, Rodney A. Badcock, and Jin Fang

Abstract—The dynamic resistance which occurs when a superconductor carrying DC current is exposed to alternating magnetic field plays an important role in HTS applications such as flux pumps and rotating machines. We report experimental results on dynamic resistance in a four-tape coated conductor stack when exposed to AC magnetic fields with different magnetic field angles (the angles between the magnetic field and normal vector component of the tape surface, θ) at 77 K. The conductors for the stack are 4 mm-wide SuperPower SC4050 wires. The field angle was varied from 0° to 120° at a resolution of 15° to study the field angle dependence of dynamic resistance on field angle as well as wire $I_c(B, \theta)$. We also varied the field frequency, the magnetic field amplitude, and the DC current level to study the dependence of dynamic resistance on these parameters. Finally, we compared the measured dynamic resistance results at perpendicular magnetic field with the analytical models for single wires. Our results show that the dynamic resistance of the stack was mainly, but not solely, determined by the perpendicular magnetic component. $I_c(B, \theta)$ influences dynamic resistance in the stack due to tilting of the crystal lattice of the superconductor layer with regard to buffer layers.

Index Terms—Dynamic resistance, HTS stack, Angle dependence.

I. INTRODUCTION

HTS coated conductors have attracted increasing interest in many HTS applications such as rotating machines, magnets, flux pumps and Superconducting Energy Storage Systems (SMES) [1]-[7]. One concern for those applications is dynamic resistance in HTS wires which occurs when the HTS wires carry DC transport current under an alternating magnetic field [8]-[9]. Dynamic resistance results in energy loss, which may cause challenges in the cryogenic system and HTS coil windings. In HTS rotating machines, HTS rotor coils carry DC current exposed to harmonic ‘ripple’ fields [1], [10]-[13]. There is no measureable dynamic resistance in the rotor windings when the amplitude of external ‘ripple’ field smaller than a threshold

magnetic field. However, as soon as the amplitude of the ‘ripple’ field becomes larger than the threshold magnetic field, a measureable dynamic resistance will occur due to the interaction between DC current and traversing flux [14]-[15]. Understanding dynamic resistance characteristics can provide useful information for HTS rotor designers. In HTS flux pumps, dynamic resistance determines the maximum pumped current and can also be utilized for current switches [16]-[17].

Other experimental and numerical reports on dynamic resistance have focused on single coated conductors under applied magnetic field with various field angles [18]-[19]. In our previous work, we measured a four-tape YBCO stack in perpendicular magnetic field, where each conductor composing the stack carries the same current [20]. In reality, in HTS applications, field angle (the angle between the magnetic field and normal vector component of the coated conductor surface, θ as defined in Fig. 1) may take any value, rather than always being zero (perpendicular magnetic field). However, there has been no reports on the dependence of dynamic resistance in HTS stacks exposed to external magnetic field with various orientations.

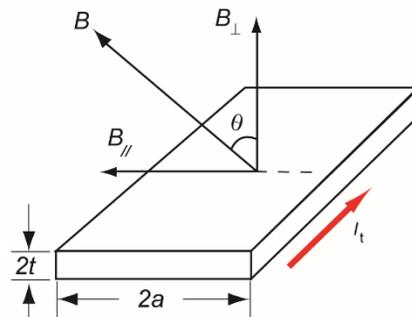


Fig. 1. Definition of field angle, θ

In this work, we present the dependence of dynamic resistance on the applied field angle in a four-tape YBCO stack.

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The field angle was varied from 0° to 120° at a resolution of 15° to investigate the field angle dependence of the dynamic resistance. The DC current level, the magnetic field amplitude, and the frequency of the magnetic field have been varied to study the influence of these parameters on the dynamic resistance. The measured results in perpendicular magnetic field were compared with the analytical models for single wires.

II. EXPERIMENTAL METHOD

Fig. 2 shows the schematics of our experimental set up to measure dynamic resistance in the four-tape stack under external magnetic field with various field angles [20]. An AC magnet composed of two vertical racetrack coils was used to produce AC magnetic field amplitudes up to 100 mT [20]. Two resonant frequencies of the magnetic field were chosen by adjusting the value of capacitance of a serially connected capacitor bank in the circuit to explore the frequency dependence of the dynamic resistance of the stack. A CROWN K2 audio amplifier was used to energize the magnet. The HTS stack was mounted on the top of the machined surface of a GFRP sample holder and then placed in the center of the magnet. The field angle, θ , was varied by rotating the sample holder inside the magnet. In this work, a 4 mm-wide SCS4050 YBCO wire with self-field I_c of 104 A manufactured by SuperPower Inc. was used to assemble the four-tape YBCO stack sample. The specifications of the wire are shown in Table 1. Kapton sheets and tapes were used as insulation between the wires composing the stack, and the vertical distance between the neighboring superconductor layers is ~ 0.295 mm. Each tape carries the same DC current by connecting the tapes in series. Hewlett Packard 6682A DC power supply was used to drive the DC sample circuit. Spiral loops attached on each conductor as shown in Fig. 2 were used to measure voltage in each conductor due to dynamic resistance [21]. The effective distance between the voltage taps is 80 mm. The DC voltage output from the voltage taps was measured using a dc Hewlett Packard 34420A nano-volt meter.

All measurements were carried out at 77 K. Both the magnet and the stack sample were immersed in the liquid nitrogen.

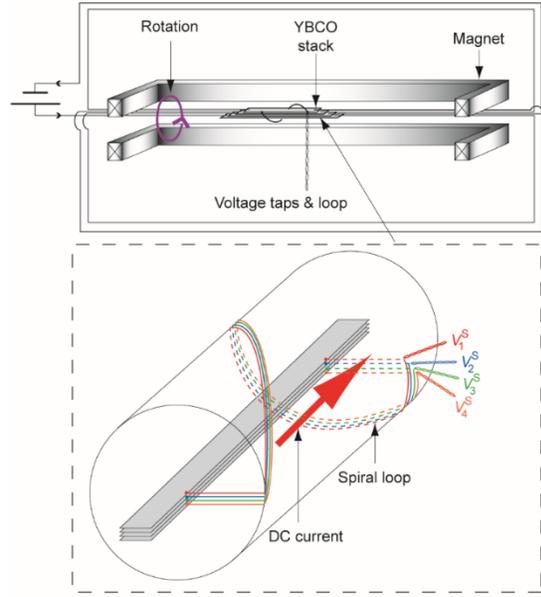


Fig. 2. Experimental set-up for stack dynamic resistance measurement and voltage tap arrangement of for the stack

TABLE 1. SPECIFICATIONS OF THE WIRE FOR STACK SAMPLE

Manufacture	SuperPower (SCS44050)
Self-field critical current (A)	104.0
Sample width (mm)	4.0
Thickness of Cu stabilizer layer each side (μm)	20
Thickness of superconductor layer (μm)	1.0
Thickness of substrate (μm)	50.0
Distance between neighboring superconductor layers (μm)	295.0

III. EXPERIMENTAL RESULTS

A. I_c measurement of the wires before and after assembling the stack

The I_c values of the four wires before and after assembling the stack are shown in Table 2. The I_c values of the wires became smaller in the stack environment than self-field I_c values of the single wires due to the superposition of the magnetic fields generated by other tapes [20].

The I_c values of the two outer tapes (T1 and T4) are smaller than those of the two inner tapes (T2 and T3) under stack environment due to the outer tapes experiencing larger magnetic fields than the inner tapes [22].

TABLE 2. I_c VALUE OF WIRE BEFORE AND AFTER ASSEMBLING

	T1	T2	T3	T4	Average
Before assembling (A)	105.7	105.3	104.6	104.6	105.1
After assembling (A)	81.6	85.4	85.3	83.9	84

B. R_{dyn} values in the stack under perpendicular magnetic field

The dynamic resistance for a single coated conductor shown in Fig. 1 under perpendicular magnetic field per unit length per cycle, R_{dyn} , can be calculated by [23]-[24],

$$\frac{R_{dyn}}{fl} = \frac{4a}{I_{c0}} (B_a - B_{th}) \quad (1)$$

where, I_{c0} is the self-field critical current of the conductor, a is the half-width of the conductor, l is the distance between the two voltage taps, B_a is the amplitude of the applied magnetic field, and f is the frequency of the applied magnetic field. B_{th} is the threshold magnetic field, which is given by,

$$B_{th} = B_{ep}(1 - i) = B_{ep} \left(1 - \frac{I_t}{I_{c0}}\right) \quad (2)$$

where i is I_t/I_{c0} and $(1 - I_t/I_{c0})$ is the filling portion of the DC current and B_{ep} is the effective penetration field [25]-[26]. In our previous work, we have derived an analytical equation for the calculation of the B_{th} value in coated conductors exposed to perpendicular magnetic fields. We have shown that $B_{ep} = 2.4642\mu_0 t J_{c0}/\pi$ [23]. An alternative equation of B_{th} was derived by Mikitik and Brandt, $B_{th,MB}$ [24],

$$B_{th,MB} = \frac{\mu_0 J_{c0} t}{2\pi} \left[\frac{1}{i} \ln \left(\frac{1+i}{1-i} \right) + \ln \left(\frac{1-i^2}{4i^2} \right) \right] \quad (3)$$

Previous work has shown that for a single coated conductor, Eq. (2) provides a better agreement with the experimentally derived B_{th} values for $i \geq 0.2$ than the values from Eq. (3). However, for $i < 0.1$, B_{th} values from Eq. (3) have better agreement with the experimentally derived B_{th} [23].

We measured dynamic resistance values of each tape in the stack at 67.89 Hz and 87.65 Hz for different DC current levels under perpendicular AC magnetic fields to investigate the frequency dependency of dynamic resistance in the stack. The measured dynamic resistance values at the two frequencies agreed with each other, and this implies the hysteretic nature of dynamic resistance in the stack [20].

Fig. 3 shows the measured R_{dyn} in each tape for 5 different current values at 67.89 Hz in perpendicular field (0°). The R_{dyn} are normalized by the sample length and magnetic field frequency, and are plotted as a function of the amplitude of the applied magnetic field. The theoretical values in a single coated conductor obtained from Eqns. (1) and (2) are also shown in the figure. The result shown in Fig. 3 are similar to the results in our 4-tape stack even though the wire I_c values used in the stacks in this work are different [20]. In the following we recap the main characteristics of the result in the stack:

- 1) The R_{dyn} values in each conductor at each current level broadly agree with one another.
- 2) The B_{th} values for the stack (which are extracted from the x -axis intercept of linear fits of the measured data) decrease with increasing DC current levels.
- 3) The B_{th} values for the stack are bigger than those in the single conductor at each DC current. This is due to the shielding effect arising from circulating currents in each conductor in the stack.
- 4) The gradient of the linear fits (dR_{dyn}/dB_a) for each tape in the stack are much larger than the single conductor.

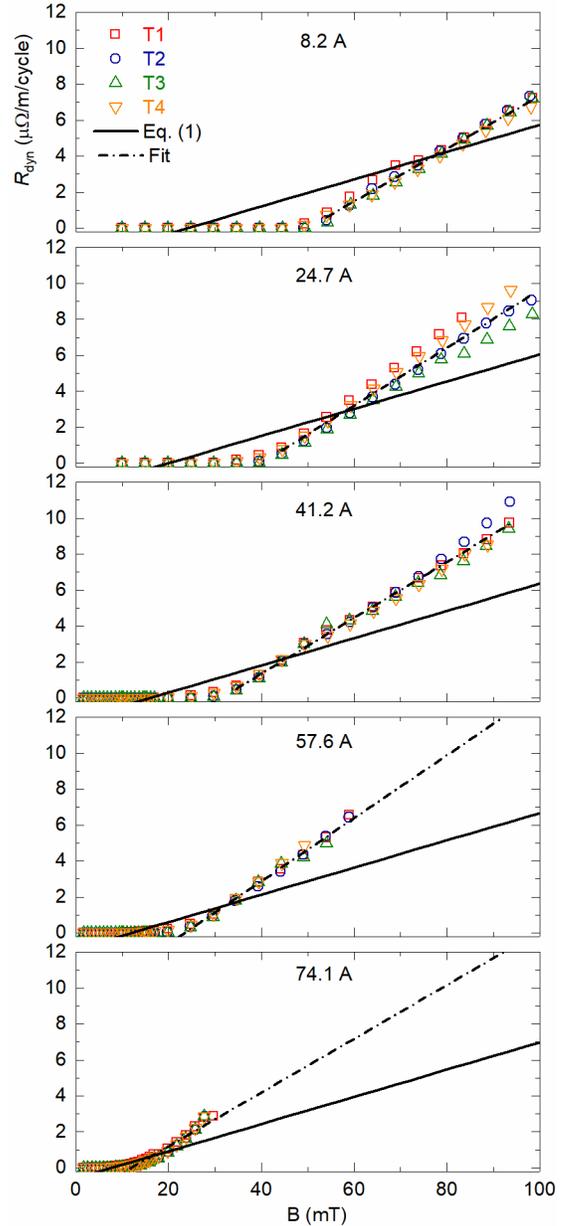


Fig. 3. Measured dynamic resistance in each wire of the YBCO stack at 5 different current for 67.89 Hz under perpendicular magnetic field compared with the calculated results in the single tape by Eq. (1).

Fig. 4 shows the comparison of the experimentally obtained B_{th} values for the stack (one is from T1 and the other is the averaged value of all four tapes), measured and analytical B_{th} values for the single conductor in perpendicular magnetic field at various different DC current values ranging from 1 A to 100 A. The analytical B_{th} values are from Eqs. (2) and (3). It is worth noting that the B_{th} values for the stack from T1 and from the averaged value of the B_{th} values in all four tapes agree from one and another. As we observed in [20] before, the B_{th} values of the stack are bigger than those for the single conductor in wide range of current values (8.2 A – 74.1 A) due to the shielding effect. The difference between the stack and the single wire increases with decreasing the DC current value because the shielding effect is more significant at small transport current.

When I_t is bigger than 8.2 A, the B_{th} of stack approximately decreases linearly with increasing the DC current. The trend is similar to single tape in the same DC current range. However, interestingly, when $I_t < 8.2$ A, the B_{th} values for the stack approach the analytical values from Eq. (3). For a single coated conductor, when I_t is very small, the B_{ep} value increases due to the increase in the shielding currents [27]. As a result, B_{th} values for the single conductor show good agreement with Eq. (3) at very small I_t . Same thing might happen in the conductors in the stack [27].

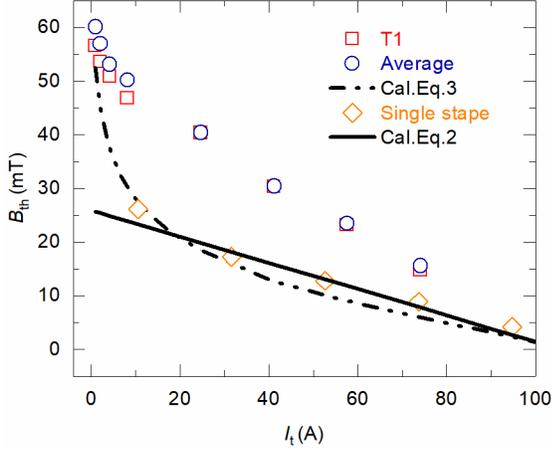


Fig. 4. Comparison of the B_{th} values of the stack and the single tape.

C. Angle dependence of the dynamic resistance

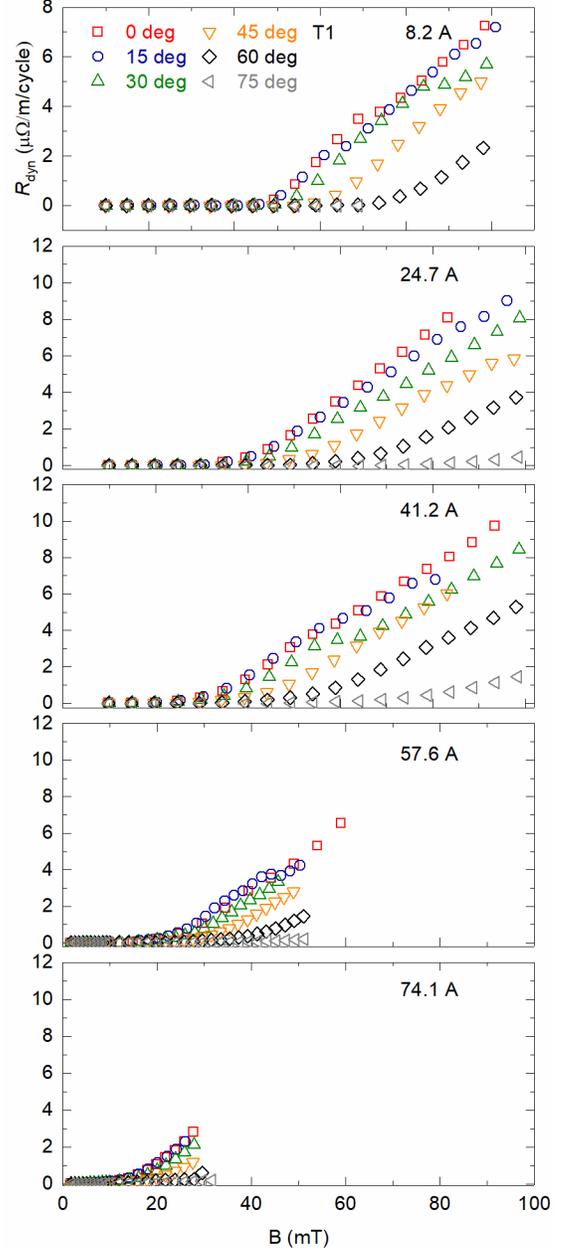


Fig. 5. Measured R_{dyn} values of T1 at various field angles plotted as a function of the applied magnetic field amplitude at 5 different currents at 67.89 Hz.

Fig. 5 shows the measured R_{dyn} values obtained from 0° to 75° at a resolution of 15° for 5 applied current values plotted as a function of the magnetic field amplitude. The results for all the tapes are similar in all cases, therefore we only show T1 result here. The results in the stack are similar to the results observed in our previous study for single wires, i.e. dR_{dyn}/dB_a becomes larger with decreasing θ and the B_{th} values decrease with decreasing θ [27]. No measurable values of R_{dyn} under parallel magnetic field ($\theta = 90^\circ$) at any of the transport currents were observed in this work.

In [27], it was concluded that the R_{dyn} of single coated conductor are mainly determined by the perpendicular component of the applied magnetic field, and the parallel magnetic field plays almost no role in the dynamic resistance of a coated conductor - due to the large aspect ratio (a/t) compared with the

BSCCO wires. In order to check the effect of the perpendicular component of the applied magnetic field on the R_{dyn} of the YBCO stack, we have replotted the data as a function of the perpendicular component of the applied magnetic field, $B_a \cdot \cos\theta$, as shown in Fig. 6.

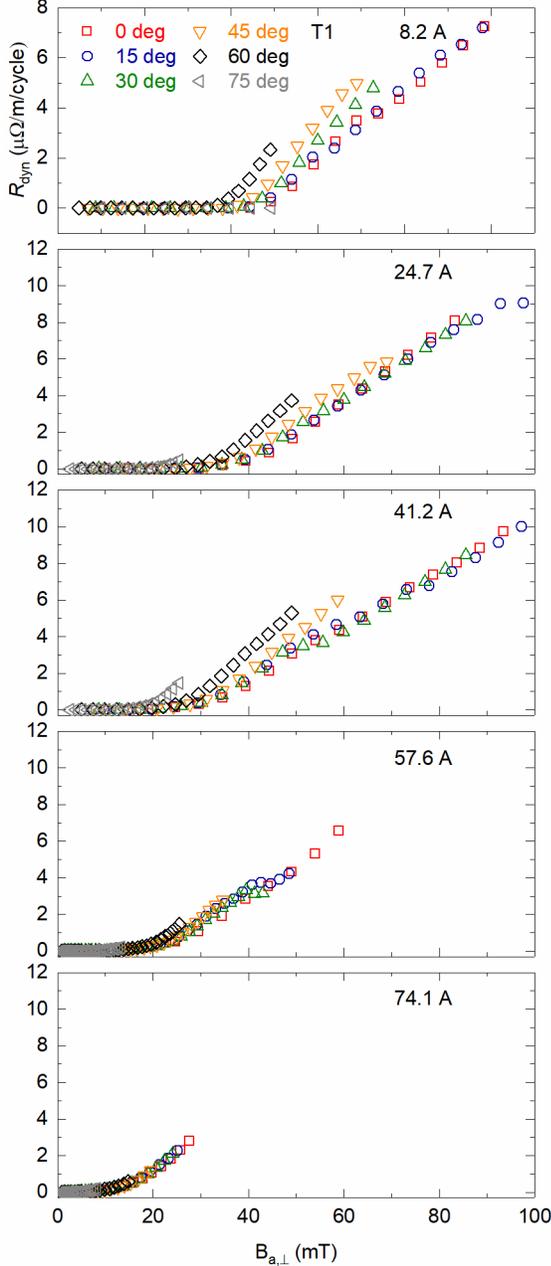


Fig. 6. Measured R_{dyn} values of T1 at various field angles plotted as a function of perpendicular magnetic field component $B_a \cdot \cos\theta$. Data obtained at 5 different currents at 67.89 Hz.

As shown in the figure, the R_{dyn} values for $\theta = 0^\circ$, 15° , and 30° broadly agree with each other for all the current values. However, the results for larger field angles deviate from the main curve. The agreement between the results for different field angles improves with increasing the DC current values. However it is clear that –in contrast to the single wire case [27]– a simple perpendicular field dependence is not sufficient to describe the angular dependence of the stack. The difference in the field angle dependence of the dynamic resistance might be due to the additive effect of off-angle field contributions from each set of

planar shielding currents flowing in the other tapes in the stack. This is a complex effect which will require more advanced modelling to fully elucidate. Additional differences with the single tape case include the relatively smaller aspect ratio of the stack than that of a single coated conductor, and the potential for imperfect planar alignment of each conductor in the stack.

D. Dynamic resistance characteristics under perpendicular magnetic field with low field amplitude

In Fig. 7, R_{dyn} values under perpendicular magnetic field at 57.6 A and 74.1 A are plotted as a function of magnetic field amplitude. The magnetic field amplitude is up to 18 mT and there was no measurable R_{dyn} when the current values were 8.2 A, 24.7 A, and 41.2 A. As shown in the figure, the R_{dyn} values of the two outer tapes (T1 and T4) are larger than those of the two inner tapes (T2 and T3), while the B_{th} values of the outer tapes are smaller than those of the inner tapes. Similar phenomenon were observed at all other angles. Qualitatively this can be explained using Eqs. (1) and (2) by replacing I_{c0} in the equations with the I_c value of each conductor in the stack arrangement. The I_c values of the outer tapes are smaller than those in the inner tapes of the stack as shown in Table 2. Therefore, at the same DC current value, the i values for the outer tapes is larger than those of the inner tapes, and this leads to smaller B_{th} values and larger R_{dyn} values. The phenomenon disappears at high magnetic field amplitudes. This is because the difference in I_c values of all the tapes becomes smaller at high magnetic field amplitudes.

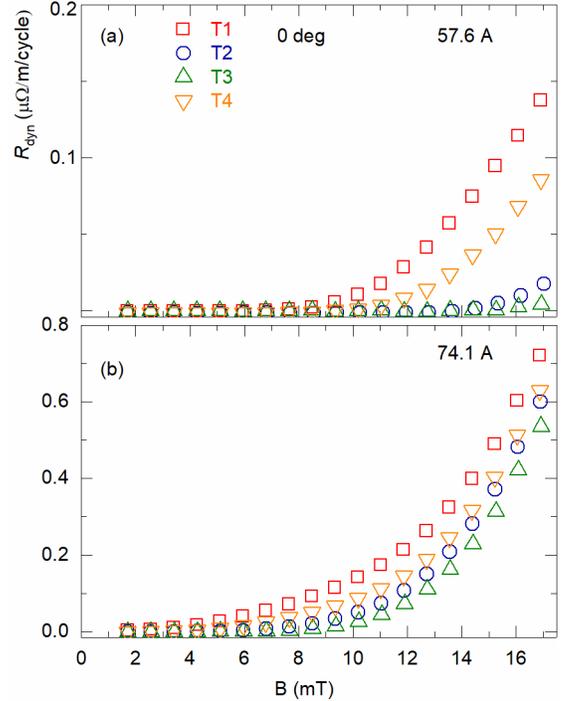


Fig. 7. Measured R_{dyn} values in each wire under perpendicular magnetic field with the magnetic field amplitude up to 18 mT at 57.6 A and 74.1 A.

E. Influence of $I_c(B, \theta)$ of sample wires on dynamic resistance

Fig. 8 shows the measured R_{dyn} values of T1 at $\theta = 60^\circ$ and 120° at 5 DC current values plotted as a function of the amplitude of the applied magnetic field. It is obvious that the measured R_{dyn} value at 60° is larger than 120° at all current values. We attribute the result to the anisotropy in the angle dependence of I_c of the coated conductor wires used in the stack [28]. Fig. 9 shows the measured I_c values of a coated conductor, which is cut from the same source material as the conductors used in the stack in this work, at different field angles at 77 K [19]. The I_c values for $\theta = 60^\circ$ and 120° are different as shown in Fig. 9 which is impossible for an isotropic coated conductor. It is immediately apparent that R_{dyn} values will be different with different I_c values according to Eqs. (1) and (2).

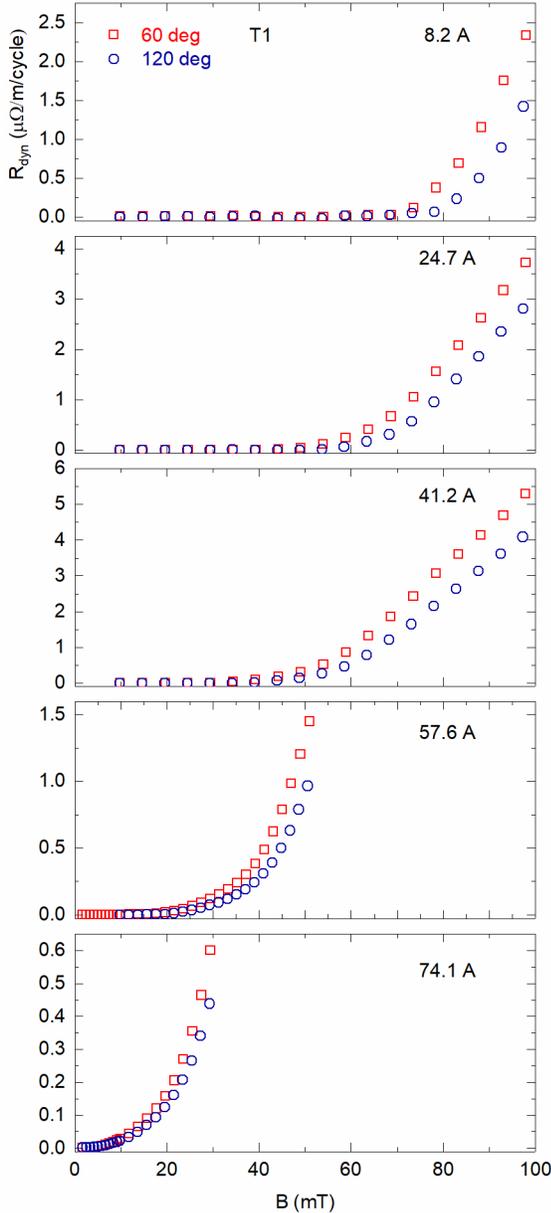


Fig. 8. Comparison of the R_{dyn} values in T1 at 60° and 120° for 67.89 Hz under 5 different current values.

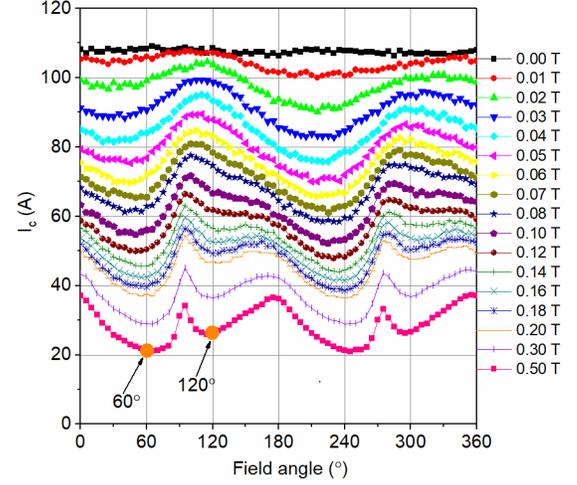


Fig. 9. Measured $I_c(B, \theta)$ values for a single coated conductor at 77 K [19].

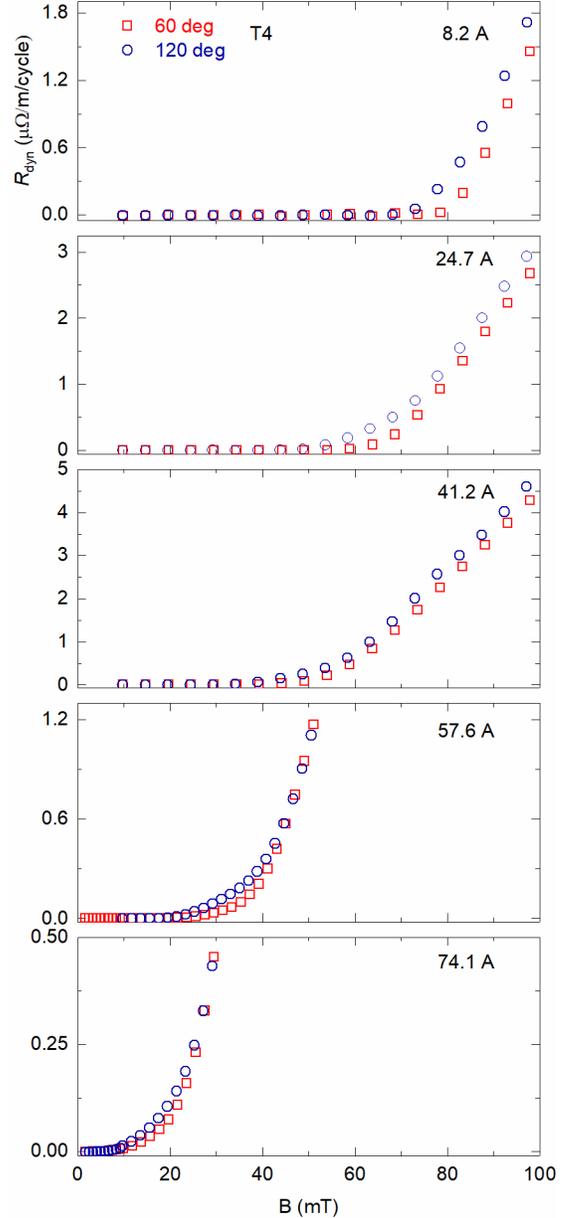


Fig. 10. Comparison of the R_{dyn} values in T4 at 60° and 120° for 67.89 Hz under 5 different current values.

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Fig. 10 shows the measured R_{dyn} values in T4 for $\theta = 60^\circ$ and 120° . Interestingly, the R_{dyn} values for $\theta = 60^\circ$ are smaller than those for $\theta = 120^\circ$ which is opposite to the results in the other three wires. We attribute the result to possible change of direction of T4 as compared to the other three wires during the sample preparation, i.e. T4 might have been accidentally rotated 180° in the same sample plane. The crystal lattice of YBCO superconductor layers is normally tilted from the substrate interface, and this results in asymmetry of the $I_c(B, \theta)$ dependence [29]. If we consider a coated conductor exposed to the external magnetic field shown in Fig. 11 (a) which is parallel to the orientation of crystal lattice of the superconductor layer. If the rotation mentioned above happens, the relative angle change between the two cases is β , as shown in Fig. 11 (b). Therefore, field angle values for the two cases (T1, T2, T3 vs T4) should be different. This gives an explanation in the result from T4. Similar results have been observed for $\theta = 75^\circ$ and 105° .

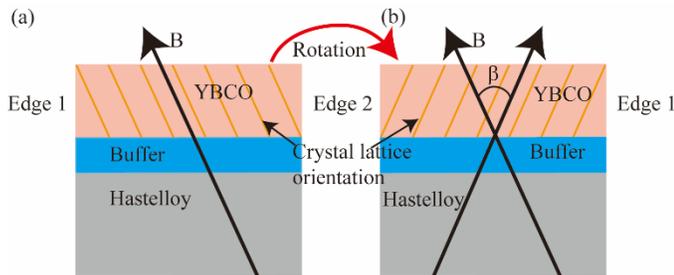


Fig. 11. Schematics of the change of relative angle between the external magnetic field and the orientation of crystal lattice in T4 due to rotation of the tape in the same plane: (a) Zero relative angle before rotation, (b) relative angle β after rotation.

IV. CONCLUSION

We have studied the dependence of dynamic resistance on the field angle of a four-tape YBCO stack while carrying DC currents exposed to AC magnetic fields with various field angles.

Shielding effect was reconfirmed through the difference in B_{th} values between the stack and a single coated conductor in this work.

At very small DC currents, i.e. when the DC current is smaller than 8.2 A, the B_{th} value for the stack agrees with theoretical value from Mikitik and Brandt.

Unlike single coated conductors, R_{dyn} values at different field angles are mainly but not solely determined by perpendicular magnetic field component possibly due to the difference in aspect ratio of the width over thickness of the stack and the single wire.

At $I_t = 57.6$ A and 74.1 A, we observed difference in R_{dyn} values for the outer two tapes and inner two tapes at low magnetic field amplitudes. This should be due to difference in I_c values in the outer and inner tapes when the tapes carry DC current.

$I_c(B, \theta)$ influences dynamic resistance in the stack. This is caused by tilting of the crystal lattice of the superconductor layer with regard to the substrate layers.

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